

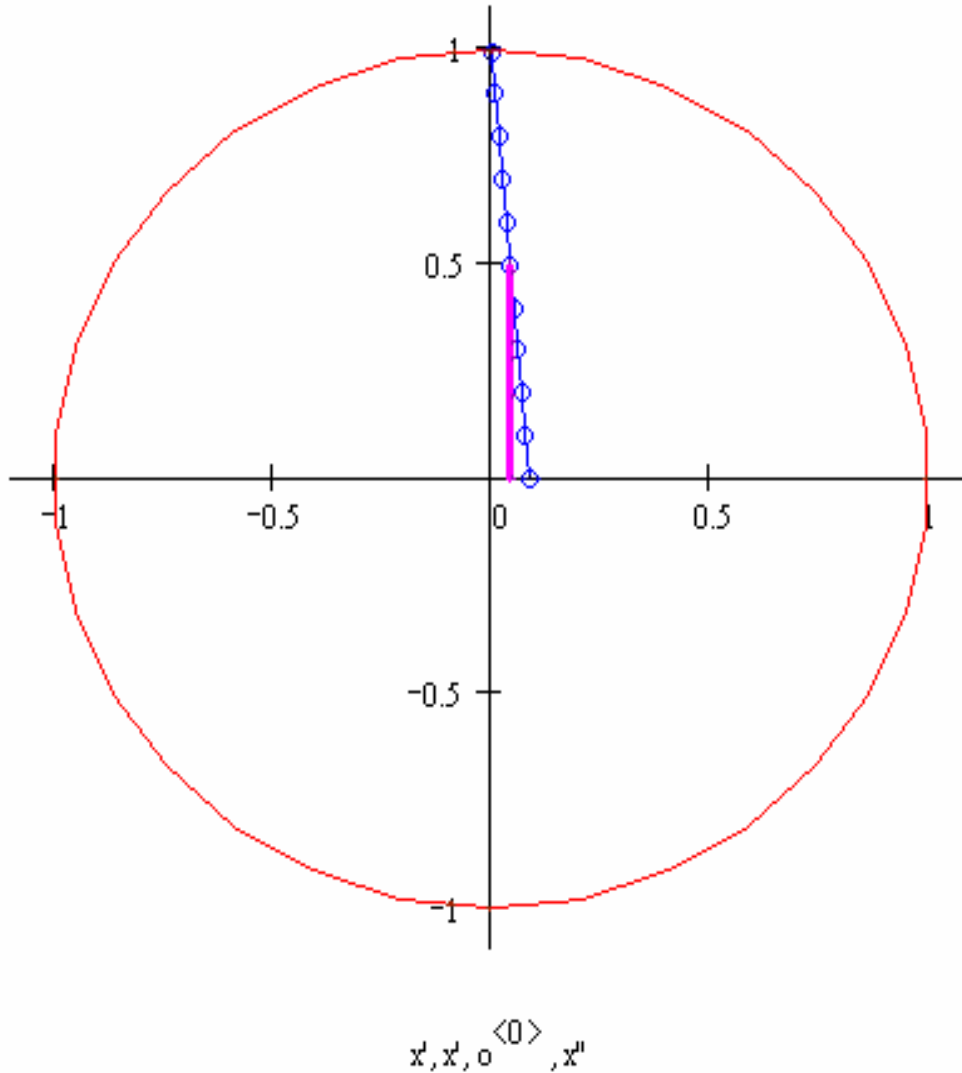
Modeling of coupled rigid-elastic multibody systems using DAE and ODE formalisms

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0.1. Example of coupled rigid-elastic system



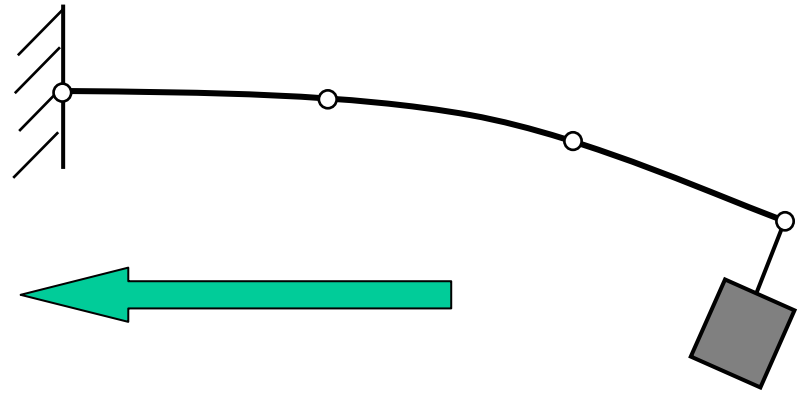
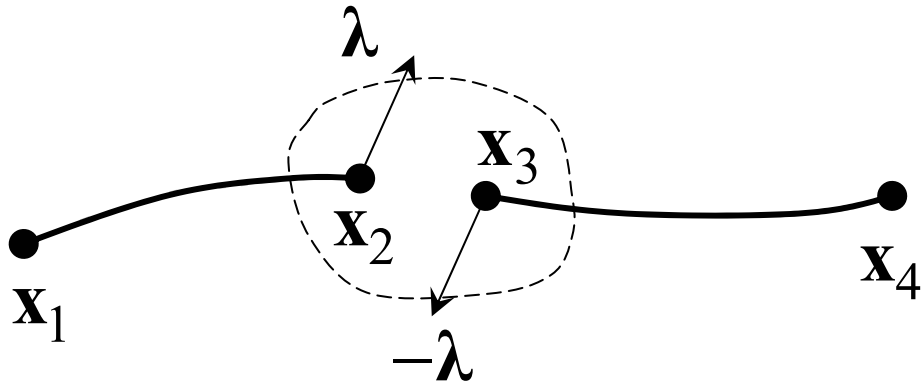
The most common way of modeling such a system employs DAE

$$\begin{cases} \mathbf{M}_1 \ddot{\mathbf{x}}_1 = \mathbf{f}_1 + (\mathbf{G}_1^T \boldsymbol{\lambda}) \\ \mathbf{M}_2 \ddot{\mathbf{x}}_2 = \mathbf{f}_2 + (\mathbf{G}_2^T \boldsymbol{\lambda}) \\ (\mathbf{g}(\mathbf{x}_1, \mathbf{x}_2) = \mathbf{0}) \end{cases}$$

$$\mathbf{G}_k = \frac{\partial \mathbf{g}}{\partial \mathbf{x}_k^T}$$

However in many cases it is possible to avoid DAE, e.g. in Finite Element Method

0.2. The Assembling Procedure in FEM



$$\begin{cases} \mathbf{M}_{11}\ddot{\mathbf{x}}_1 + \mathbf{M}_{12}\ddot{\mathbf{x}}_2 = \mathbf{f}_1 \\ \mathbf{M}_{21}\ddot{\mathbf{x}}_1 + \mathbf{M}_{22}\ddot{\mathbf{x}}_2 = \mathbf{f}_2 + (\lambda) \\ \mathbf{M}_{33}\ddot{\mathbf{x}}_3 + \mathbf{M}_{34}\ddot{\mathbf{x}}_4 = \mathbf{f}_3 - (\lambda) \\ \mathbf{M}_{43}\ddot{\mathbf{x}}_3 + \mathbf{M}_{44}\ddot{\mathbf{x}}_4 = \mathbf{f}_4 \end{cases}$$

Constraints are trivial:
 $(\mathbf{x}_2 \equiv \mathbf{x}_3 = \mathbf{x}_*)$

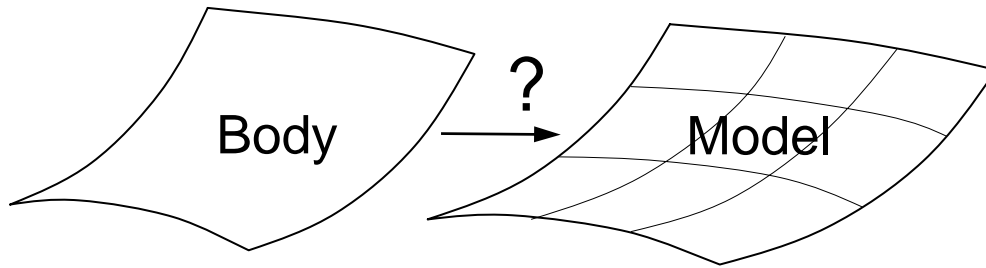
$$\begin{bmatrix} \mathbf{M}_{11} & \mathbf{M}_{12} & & \mathbf{O} \\ \mathbf{M}_{21} & \mathbf{M}_{22} & + & \mathbf{M}_{33} & \mathbf{M}_{34} \\ & & & \mathbf{M}_{43} & \mathbf{M}_{44} \\ \mathbf{O} & & & & \end{bmatrix} \begin{Bmatrix} \ddot{\mathbf{x}}_1 \\ \ddot{\mathbf{x}}_* \\ \ddot{\mathbf{x}}_4 \end{Bmatrix} = \begin{Bmatrix} \mathbf{f}_1 \\ \mathbf{f}_2 + \mathbf{f}_3 \\ \mathbf{f}_4 \end{Bmatrix}$$

Translational motion of rigid bodies is usually simulated with Cartesian coordinates.

To describe the **rotational motion** of the body we use

rotation angle φ in 2D case
or any triplet $\alpha_1, \alpha_2, \alpha_3$ of angles in 3D case.

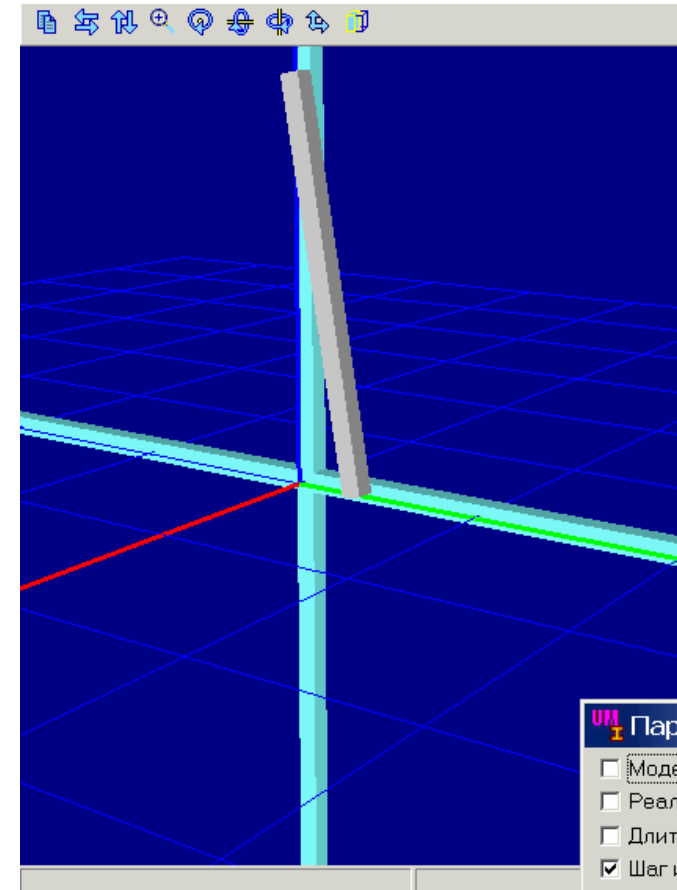
(We do not consider the Euler-Rodriguez parameters)



Consider two large-displacement approaches for simulation of elastic bodies:

1. Large rotation vector formulation
(*J. C. Simo, 1985*)
2. Absolute nodal coordinate formulation
(*A. Shabana, 1996*)

Example of large-displacement elastic system



Simulation tool is
Universal Mechanism
by Prof. *D. Pogorelov*

1.1. Large rotation vector formulation for 2D beam

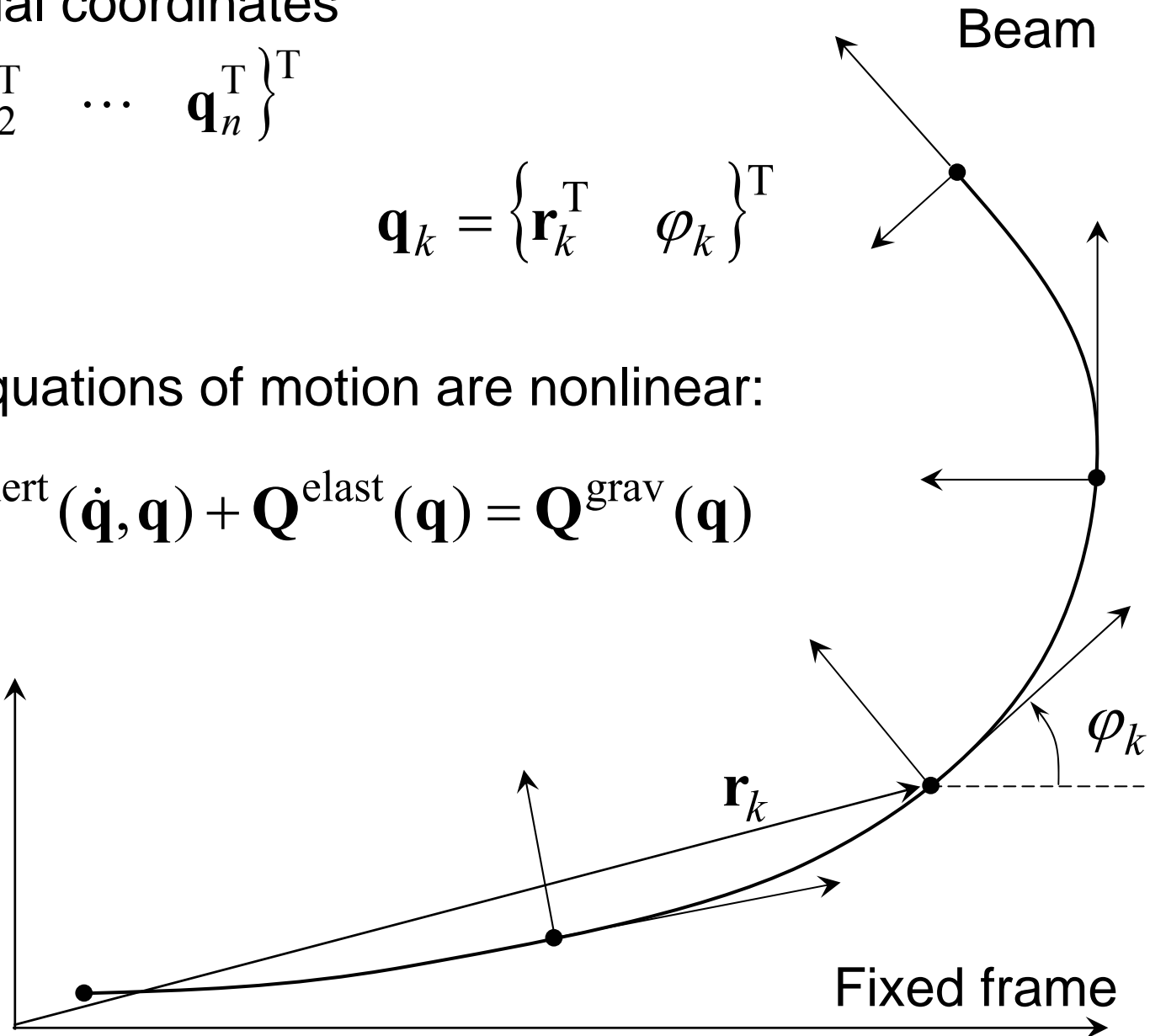
Vector of nodal coordinates

$$\mathbf{q} = \left\{ \mathbf{q}_1^T \quad \mathbf{q}_2^T \quad \cdots \quad \mathbf{q}_n^T \right\}^T$$

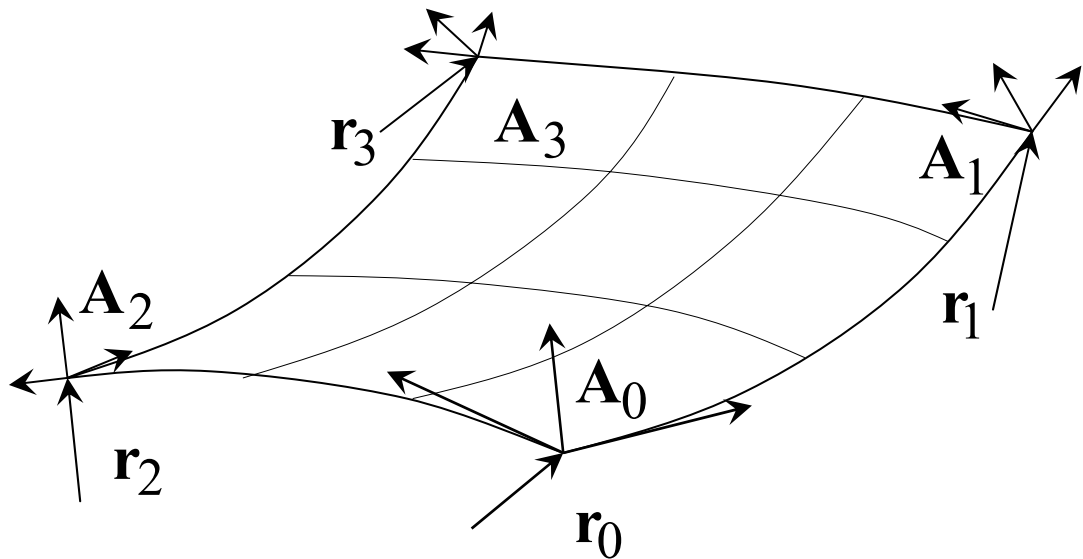
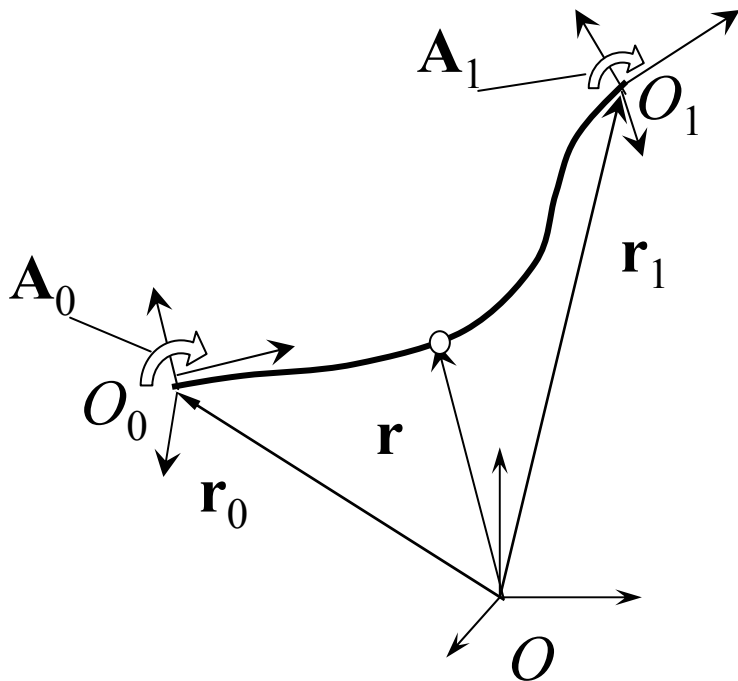
$$\mathbf{q}_k = \left\{ \mathbf{r}_k^T \quad \varphi_k \right\}^T$$

All terms of equations of motion are nonlinear:

$$\mathbf{M}(\mathbf{q}) \ddot{\mathbf{q}} + \mathbf{Q}^{\text{inert}}(\dot{\mathbf{q}}, \mathbf{q}) + \mathbf{Q}^{\text{elast}}(\mathbf{q}) = \mathbf{Q}^{\text{grav}}(\mathbf{q})$$



1.2. 3D beam and plate FEs using large rotations



Coordinates

$$\mathbf{q} = \left\{ \mathbf{q}_0^T \quad \mathbf{q}_1^T \right\}^T$$

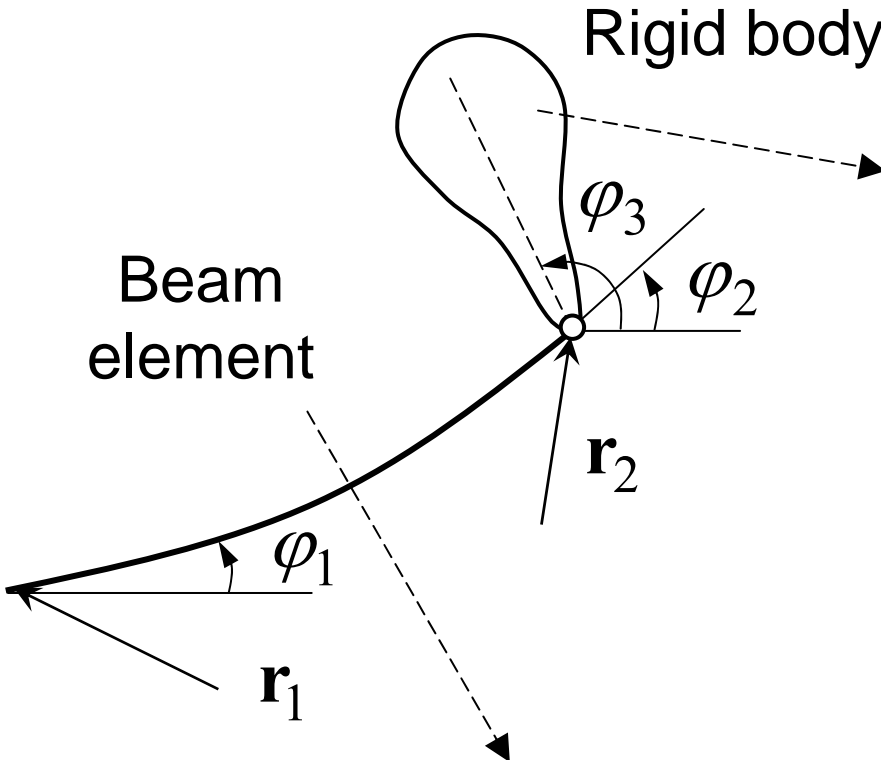
$$\mathbf{q} = \left\{ \mathbf{q}_0^T \quad \mathbf{q}_1^T \quad \mathbf{q}_2^T \quad \mathbf{q}_3^T \right\}^T$$

$$\mathbf{q}_k = \left\{ \begin{array}{c} \mathbf{r}_k \\ \boldsymbol{\varphi}(\mathbf{A}_k) \end{array} \right\}$$

A triplet of orientation angles

$$\boldsymbol{\varphi}(\mathbf{A}_k) = \left\{ \alpha_{1k} \quad \alpha_{2k} \quad \alpha_{3k} \right\}^T$$

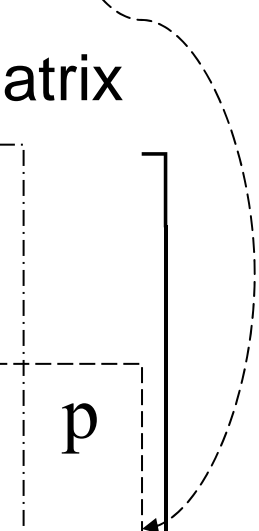
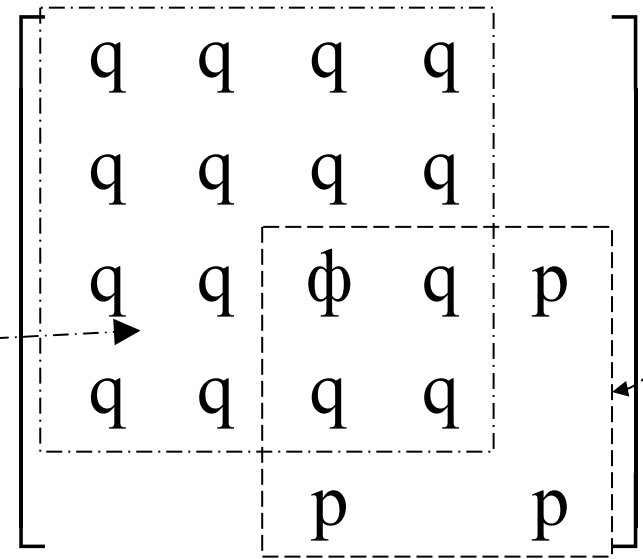
1.3. Revolute joint between beam and rigid body



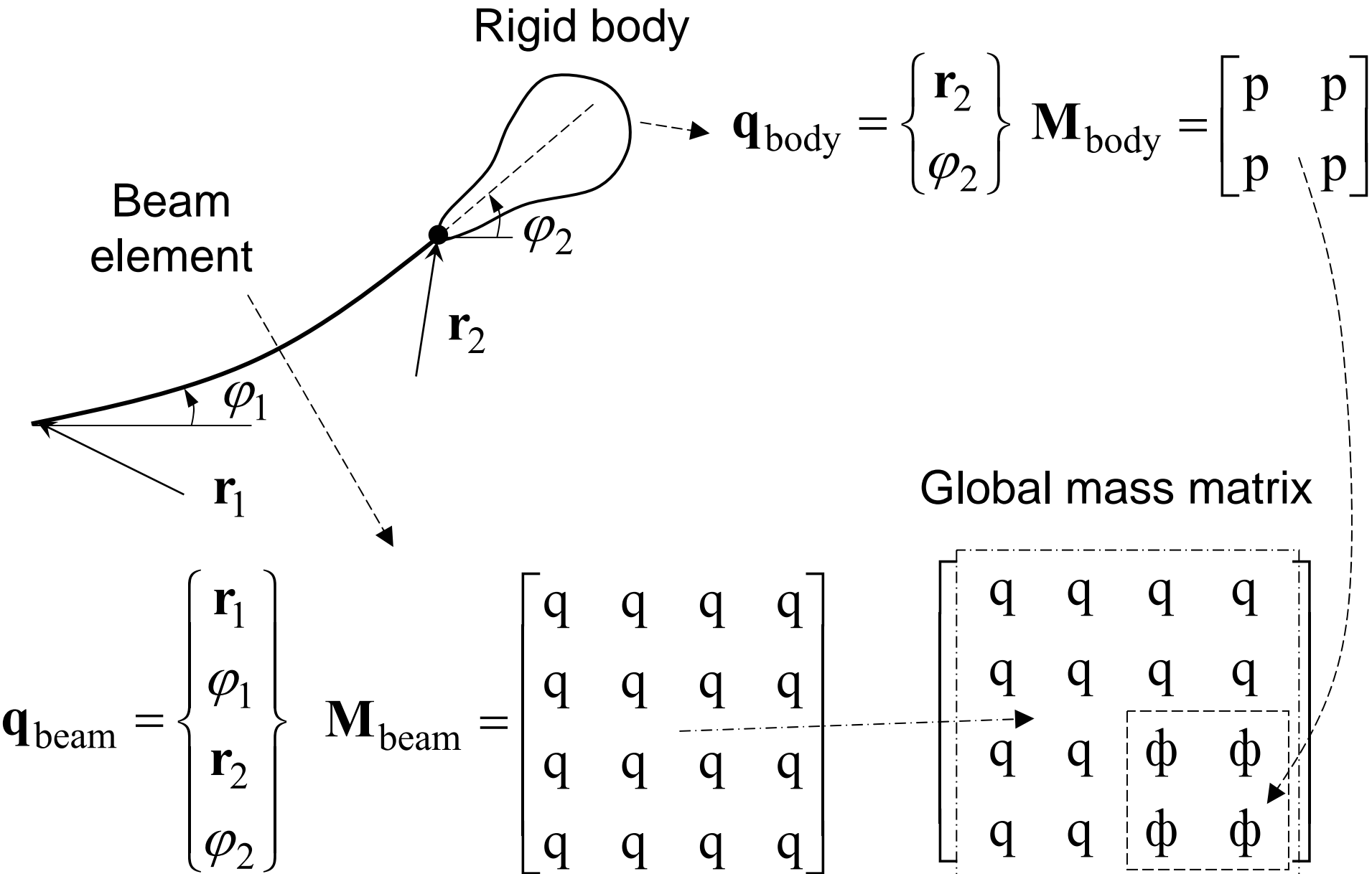
$$\mathbf{q}_{\text{body}} = \begin{Bmatrix} \mathbf{r}_2 \\ \varphi_3 \end{Bmatrix} \quad \mathbf{M}_{\text{body}} = \begin{bmatrix} p & p \\ p & p \end{bmatrix}$$

$$\mathbf{q}_{\text{beam}} = \begin{Bmatrix} \mathbf{r}_1 \\ \varphi_1 \\ \mathbf{r}_2 \\ \varphi_2 \end{Bmatrix} \quad \mathbf{M}_{\text{beam}} = \begin{bmatrix} q & q & q & q \\ q & q & q & q \\ q & q & q & q \\ q & q & q & q \end{bmatrix}$$

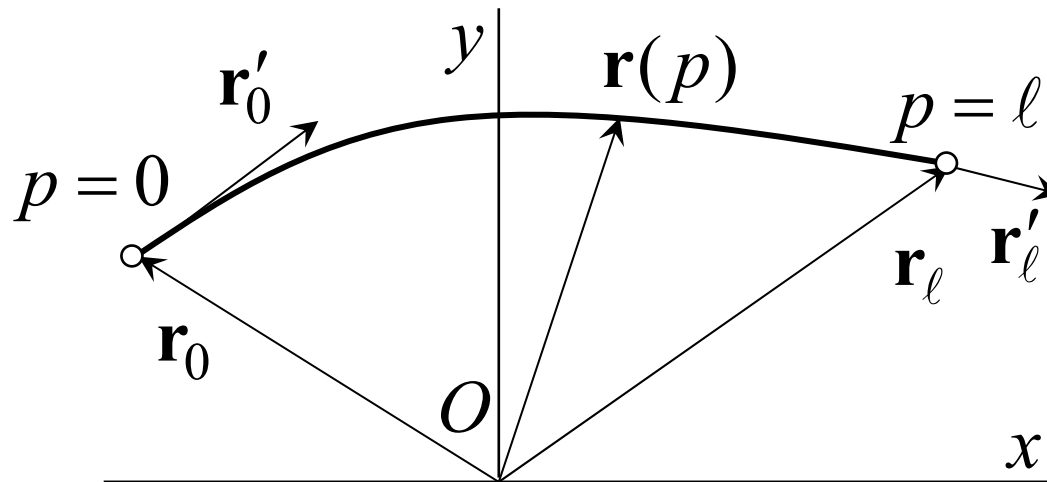
Global mass matrix



1.4. Fixed connection between beam and rigid body



2. Absolute nodal coordinate formulation (ANCF)



Radius vector
 $\mathbf{r}(p, \mathbf{q}) = \mathbf{S}(p) \mathbf{q}$

Equations of motion
 (in the large displacement case)

$$\mathbf{M} \ddot{\mathbf{q}} + \mathbf{Q}^{\text{elast}}(\mathbf{q}) = \mathbf{Q}^{\text{grav}}$$

Constant mass matrix

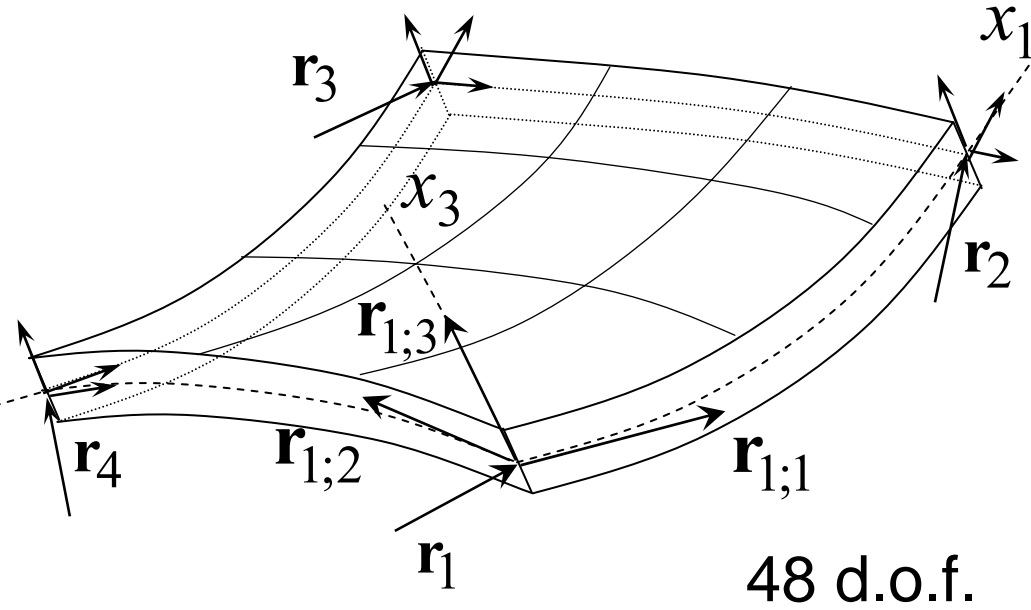
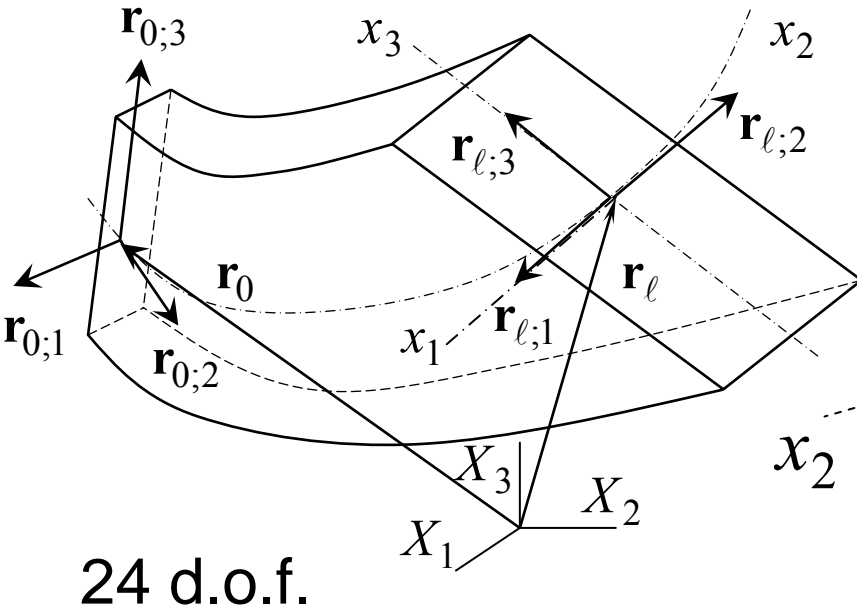
The only nonlinear term

Constant gravity
 applied forces

2.2. Thick beam/plate/shell elements by Shabana et al.

A. Shabana, R. Yakoub (2001)

A. Mikkola, A. Shabana (2001)



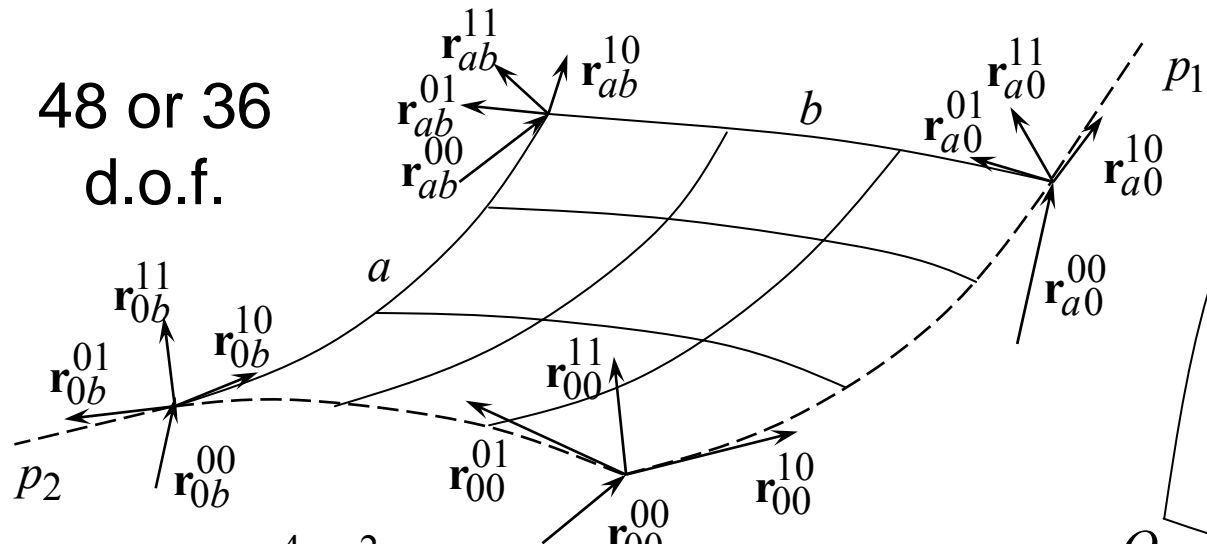
$$\mathbf{r}_{\alpha;k} = \frac{\partial \mathbf{r}_\alpha}{\partial x_k}$$

$$\mathbf{r} = \mathbf{S}(x_1, x_2, x_3) \mathbf{q}$$

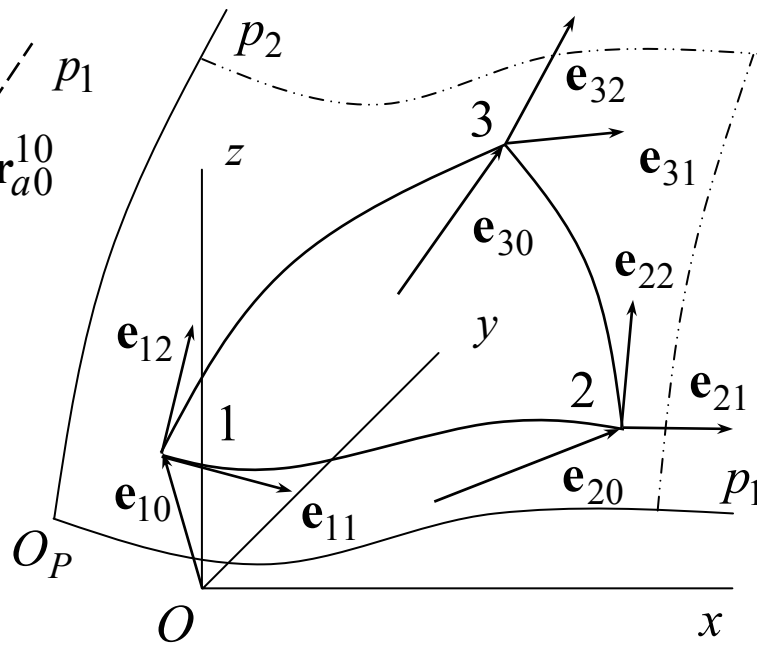
The elements represent **thick** beams / plates / shells

2.3. Thin elements by DMITROCHENKO & POGORELOV (2002)

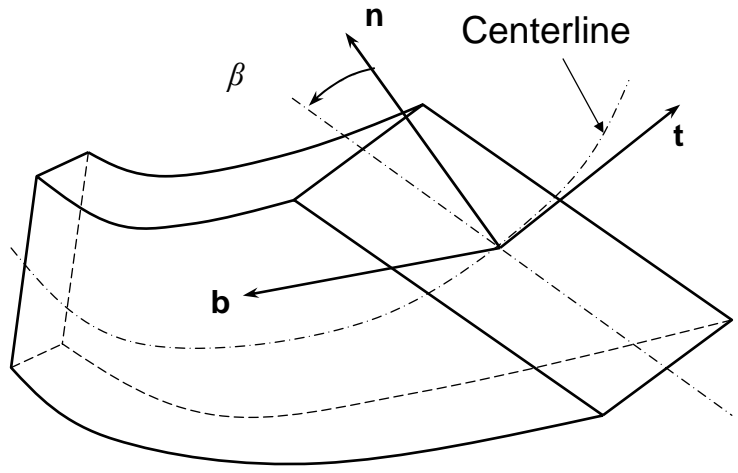
48 or 36
d.o.f.



$$\mathbf{r} = \sum_{m=1}^4 \sum_{n=0}^2 S_{mn}^{12}(p_1, p_2) \mathbf{e}_{mn}$$



27 d.o.f.

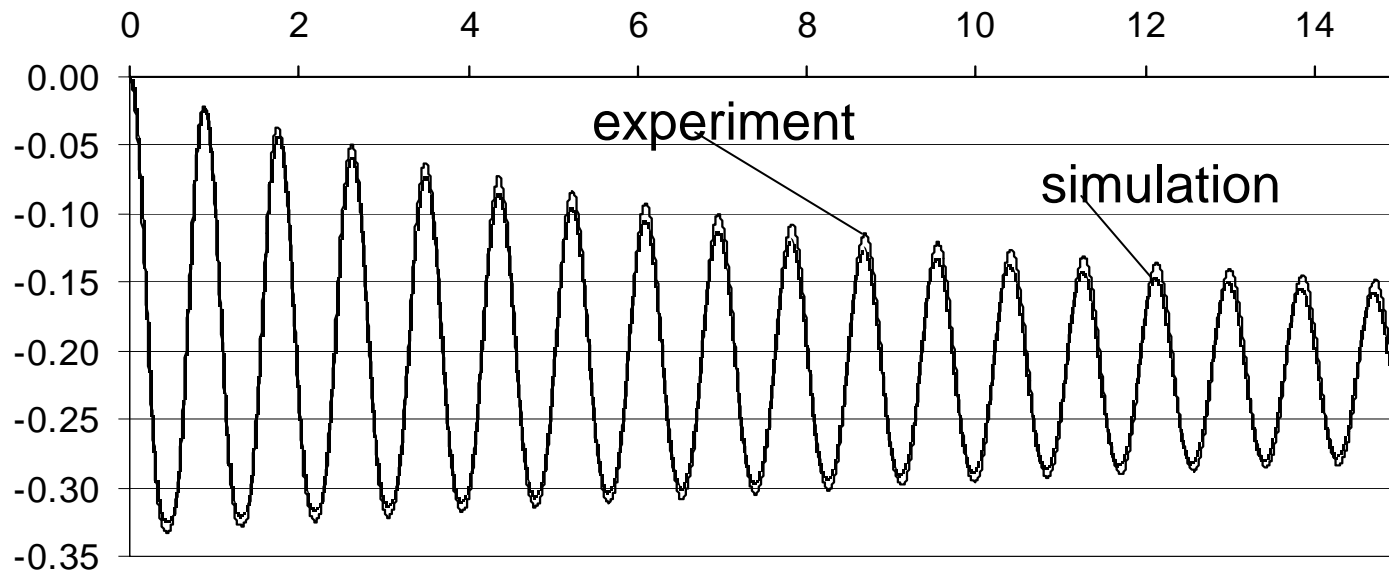
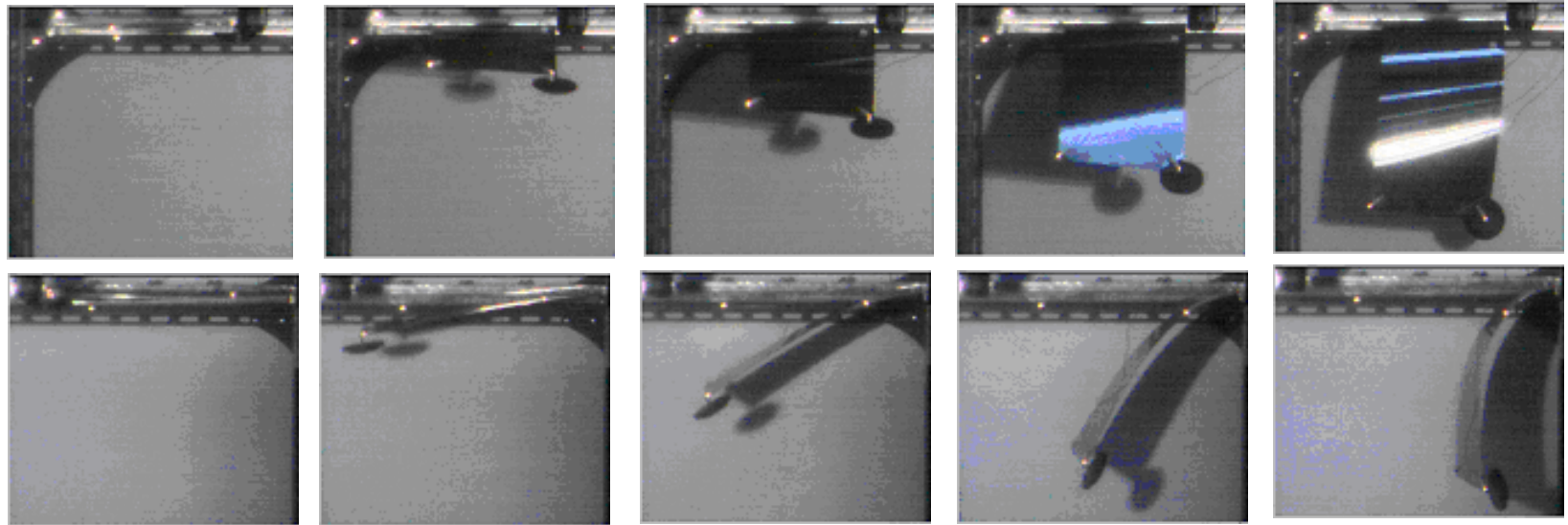


14 d.o.f.

$$\mathbf{r} = \sum_{m=1}^3 \sum_{n=0}^2 S_{mn}^L(L_1, L_2, L_3) \mathbf{e}_{mn}$$

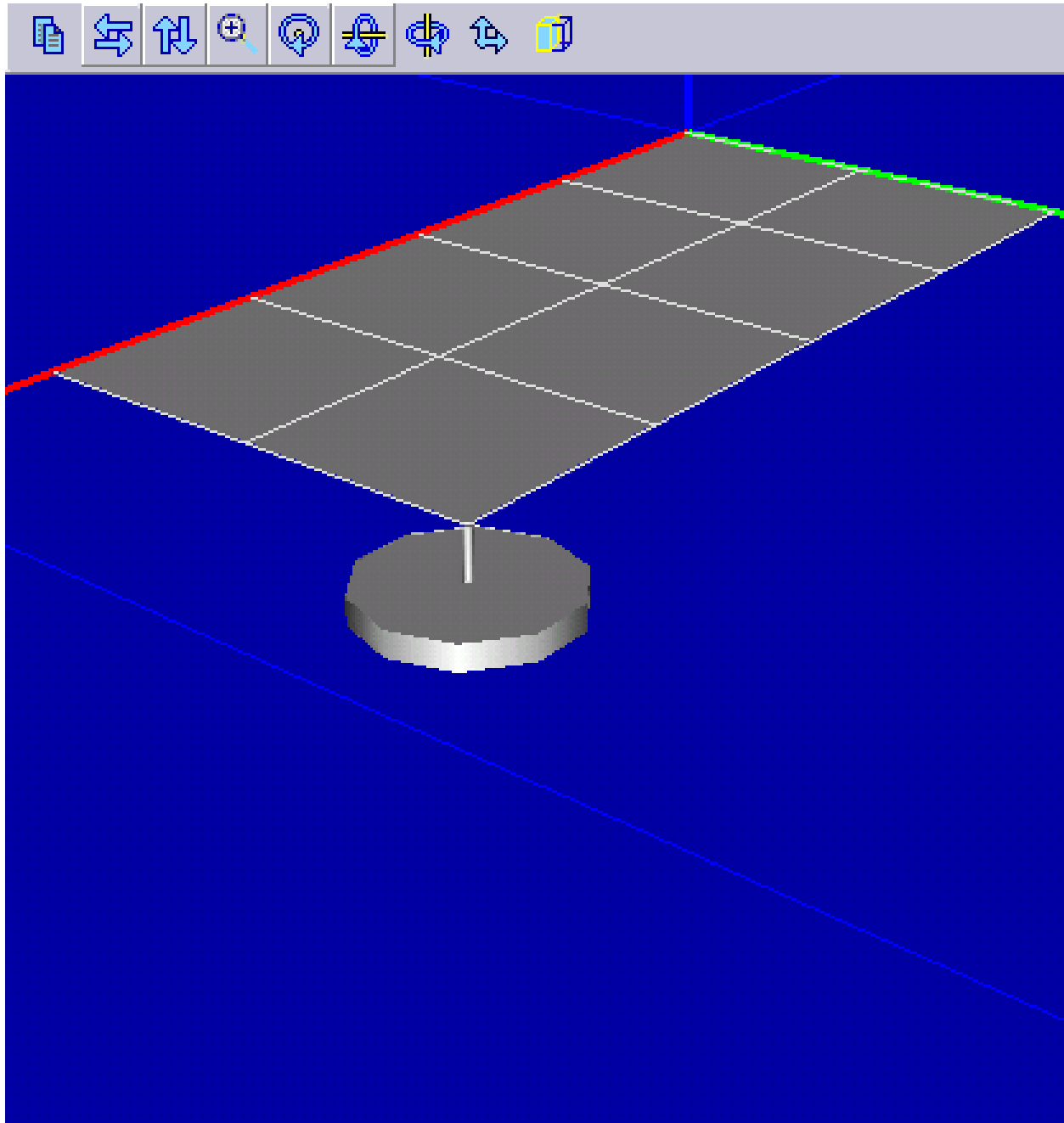
The elements represent **thin** beams and plates

Experiments were done at Pusan National University (Korea)
jointly with Prof. Wan-Suk Yoo (CAE Lab)

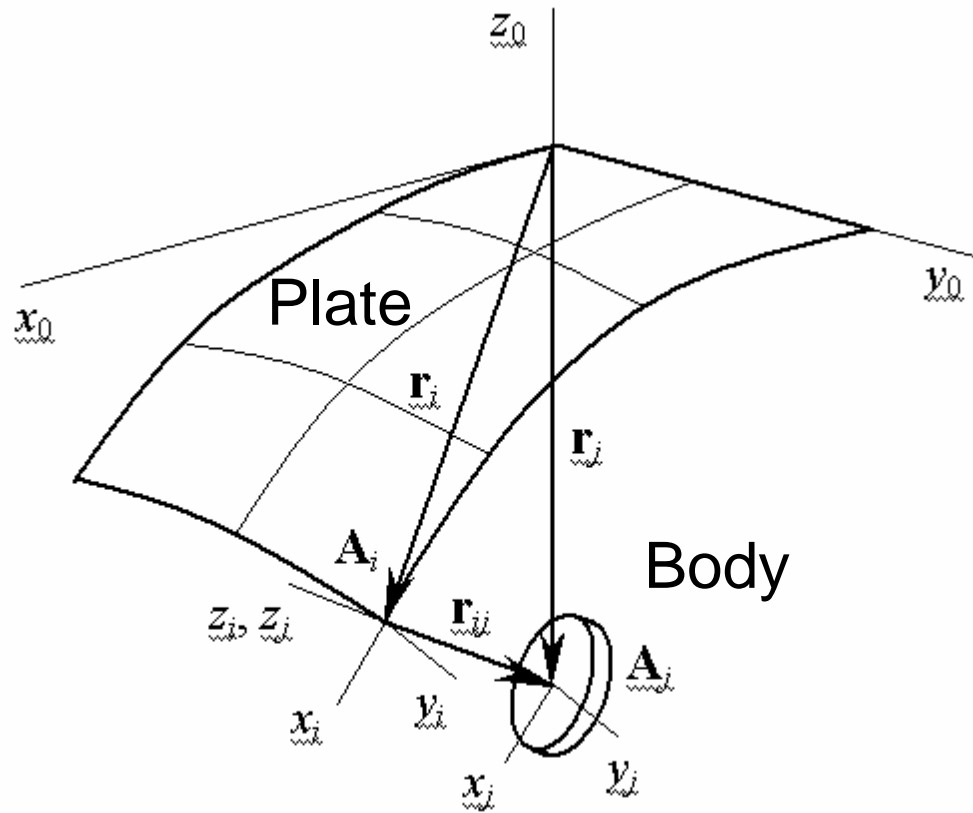
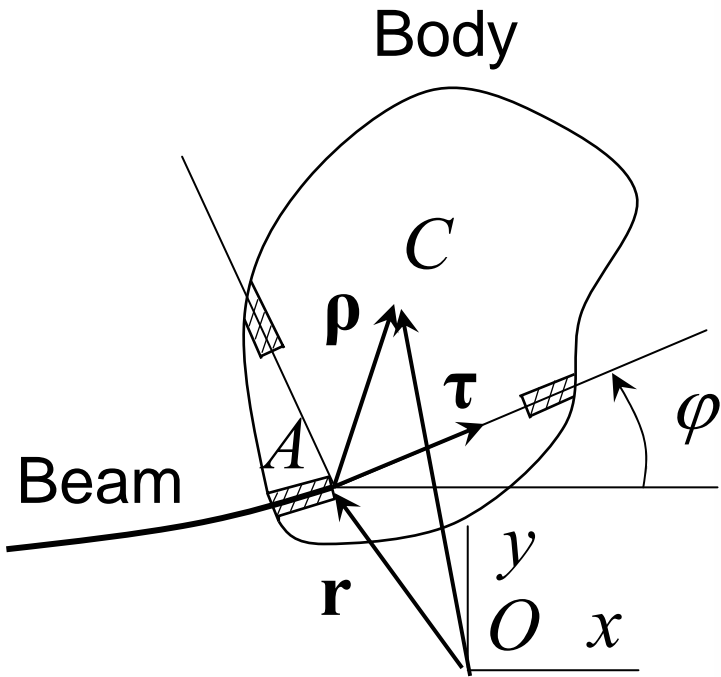


2.5. Animation of a clamped plate with a weight

Euromech 452
Colloquium 2004



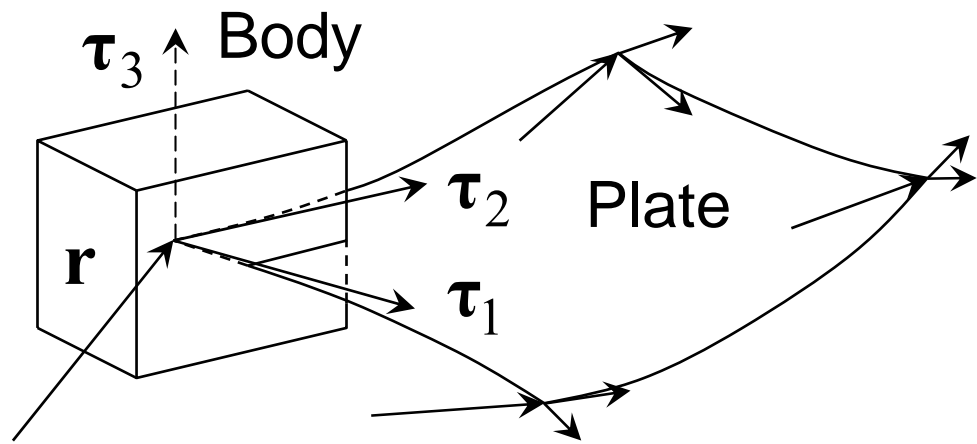
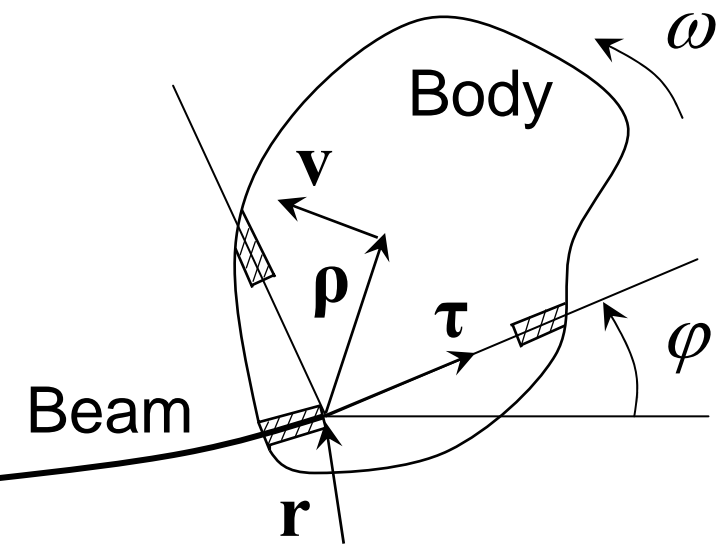
Simulation tool is
Universal Mechanism
by Prof. *D. Pogorelov*



$$\begin{cases} \mathbf{M}_i \ddot{\mathbf{q}}_i = \mathbf{Q}_i + \mathbf{G}_i^T \boldsymbol{\lambda} \\ \mathbf{M}_j \ddot{\mathbf{q}}_j = \mathbf{Q}_j + \mathbf{G}_j^T \boldsymbol{\lambda} \\ \left\{ \begin{array}{l} \mathbf{r} + \boldsymbol{\rho} - \mathbf{r}_C \\ \tan \varphi - (\tau_y / \tau_x) \end{array} \right\} = \mathbf{0} \end{cases}$$

$$\begin{cases} \mathbf{M}_i \ddot{\mathbf{q}}_i = \mathbf{Q}_i + \mathbf{G}_i^T \boldsymbol{\lambda} \\ \mathbf{M}_j \ddot{\mathbf{q}}_j = \mathbf{Q}_j + \mathbf{G}_j^T \boldsymbol{\lambda} \\ \left\{ \begin{array}{l} \mathbf{r}_i + \mathbf{r}_{ij} - \mathbf{r}_j \\ \sum_{k=1}^3 \tilde{\mathbf{u}}_k \left[\mathbf{A}_i \mathbf{A}_{ij} \mathbf{A}_j^{-1} \right] \mathbf{u}_k \end{array} \right\} = \mathbf{0} \end{cases}$$

2.7. Elimination of the algebraic part from DAE



$\mathbf{x} = \{\mathbf{r}^T \ \boldsymbol{\tau}^T\}^T$ – Coordinates of body – $\mathbf{x} = \{\mathbf{r}^T \ \boldsymbol{\tau}_1^T \ \boldsymbol{\tau}_2^T\}^T$

$\varphi = \arctg(\tau_Y / \tau_X)$
 $\omega = \dot{\varphi}(\mathbf{x}, \dot{\mathbf{x}})$

$\mathbf{v} = \dot{\mathbf{r}} + \boldsymbol{\omega} \times \boldsymbol{\rho}$
 \downarrow
 $\mathbf{v} = \boldsymbol{\Phi} \dot{\mathbf{x}}$
 $\mathbf{a} = \boldsymbol{\Phi} \ddot{\mathbf{x}} + \dot{\boldsymbol{\Phi}} \dot{\mathbf{x}}$

$\mathbf{A} \approx [\boldsymbol{\tau}_1 \ \boldsymbol{\tau}_2 \ \boldsymbol{\tau}_3]$
 $\tilde{\boldsymbol{\omega}} = \dot{\mathbf{A}} \mathbf{A}^T$

Virtual work principle

$\int_V \delta \mathbf{r}^T \mu (\mathbf{a} - \mathbf{g}) dV = 0$

Equations of motion (ODE)

$\mathbf{M}(\mathbf{x}) \ddot{\mathbf{x}} + \mathbf{Q}^{\text{inert}}(\dot{\mathbf{x}}, \mathbf{x}) = \mathbf{Q}^{\text{grav}}(\mathbf{x})$

Coupled rigid and elastic bodies can be modeled as a constraint-free problem.

We considered two finite-element large-displacement formulations to simulate elastic bodies:

- 1) **large rotation vector formulation** by J. Simo and
- 2) **absolute nodal coordinate formulation** by A. Shabana.

We considered two cases of attaching a rigid body to elastic one:

- 1) without restrictions for relative angular orientation (**revolute joint** in 2D case, **spherical joint** in 3D case)
- 2) without relative degrees of freedom (**the bodies are clamped**).

It can be shown that **this technique can be generalized** for any kind of joints between elastic and rigid bodies.

Thank you for your kind attention.